Outline	Black Hole Accretion	Cosmic Acceleration	Motive	Construction of Problem	Solutions	Conclusions

MCG Accretion : Limiting Value of η/s

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August 9, 2014



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I want to thank **Prof. Banibrata Mukhopadhyay**, Department of Physics, IISc for the core idea and fruitful discussions.

1 Black Hole Accretion

- Accretion : What is?
- Accretion : Where might be found?
- Accretion : Primary Modelling
- 2 Cosmic Acceleration
 - Cosmic Acceleration : Modified Chaplygin Gas
 - MCG around Supermassive Black Holes
- 3 Motive
 - MCG Accretion & $\frac{\eta}{s}$ Problem
- 4 Construction of Problem
 - Diff. Assumptions of Constructing Accretion Problem

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- Viscosity : Shakura Sunyaev Prescription
- 5 Solutions
 - MCG Accretion : Stronger Outflowing Winds
 - Solution : $\frac{\eta}{s}$

6 Conclusions



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- Gravity causes material in the disc to spiral inward towards the central body. Rayleigh Stability Criterion: $\frac{\partial (r^2 \omega)}{\partial r} > 0.$ Gravitational and frictional forces compress and raise the temperature of the material causing the emission of electromagnetic radiation.

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Accretion : Where might be found?

The most spectacular accretion discs found in nature are those of active galactic nuclei and of quasars, which are believed to be massive black holes at the center of galaxies.

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The large luminosity of quasars is believed to be a result of gas being accreted by supermassive black holes. This process can convert about 10% of the mass of an object into energy as compared to around 0.5% for nuclear fusion processes.

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The point L1 : Inner Lagrangian Point, The wind processed by the optical companion when fills the surrounding of it and starts to flow towards the X ray Companion through the L1. This is called Roche Lobe over flow.



Within the framework of Einstein's gravity, due to this present accelerating phase, it is reasonable to believe that DE is the dominating part (74.5% of the energy content in the observable universe)¹ of the total energy of the universe.

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- A Dark energy candidate MCG has EoS as ²

$$\boldsymbol{p} = \alpha \rho - \frac{\beta}{\rho^n} \tag{1}$$

It gives the cosmological evolution from an initial radiation era (with $\alpha = \frac{1}{3}$) to (asymptotically) the ΛCDM era (where fluid behaves as cosmological constant).

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- At present, it is widely believed that the hierarchical clustering of cold dark matter is the origin of structures in the universe. Consequently, the angular momentum of dark matter halos and eventually the rotation of galaxies is thought to be produced by gravitational tidal torque. It has been suggested that the halos obtain their spins through the cumulative acquisition of angular momentum from satellite accretion. DE that has accreted on a galaxy would be similarly torqued by such tidal interactions.

³Gamow, G. :- *Phys. Rev.* 86 251,(1952). Ritabrata Biswas, IIEST Shibpur, Howrah. TCGC-2014 MCG Accretion : Limiting Value of *n/s*



It is a question what will be the fate of the feeding process of the BH, i.e., the accretion disc.

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- Is there any natural site revealing an $\frac{\eta}{s}$ close to above lower bound?



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CENtrifugal pressure supported BOundary Layer acts as the effective boundary layer of the BH system which, like stellar surface, could produce outflowing winds.

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The Angular Momentum Gradient

$$u \frac{d\lambda}{dx} = \frac{2\sqrt{F\alpha_s}}{x^{\frac{1}{2}}c_s} \left[\frac{u^2 - c_s^2 + (n+1)\alpha_s}{n} - x(c_s^2 + u^2) \frac{1}{n+1} \frac{c_s}{c_s^2 - \alpha_s} \frac{dc_s}{dx} + xu \frac{du}{dx} \right]$$

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 Relativistic flow must transit to supersonic phase from subsonic via sonic point/ critical point.

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- We will fix the critical point distance, i.e., we will provide the energy of the accreting object. N = D = 0 relation with particular r_c will give us the radial velocity and the critical point sound speed. We will fix λ_c manually.

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The Angular Momentum Gradient

$$u \frac{d\lambda}{dx} = \frac{2\sqrt{F\alpha_s}}{x^{\frac{1}{2}}c_s} \left[\frac{u^2 - c_s^2 + (n+1)\alpha_s}{n} - x(c_s^2 + u^2) \frac{1}{n+1} \frac{c_s}{c_s^2 - \alpha_s} \frac{dc_s}{dx} + xu \frac{du}{dx} \right]$$

- Relativistic flow must transit to supersonic phase from subsonic via sonic point/ critical point.At the critical point the denominator of the velocity gradient becomes zero. To make a continuous flow we should have numerator zero as well and to use L'Hospital's rule to obtain the sonic point velocity gradient.
- We will fix the critical point distance, i.e., we will provide the energy of the accreting object. N = D = 0 relation with particular r_c will give us the radial velocity and the critical point sound speed. We will fix λ_c manually.



Solutions : Spherical Accretion



Fig. 1(a) & 1(b) represent the variation of accretion and wind speeds as functions of radial coordinate for $\lambda = 0$, j = 0. The solid line represents the accretion speed whereas the rest is for wind speed. It should be marked that the absolute value of the velocities are been plotted here.



Solutions : Disc Accretion, Non Rotating BH



Fig. 2(a) & 2(b) represent the variation of accretion and wind speeds in disk flows as functions of radial coordinate for j = 0. The solid line represents the accretion speed whereas the rest is for wind speed. It should be marked that the absolute value of the velocities are been plotted here.



Solutions : Disc Accretion, Rotating BH



Fig. 3(a) & 3(b) represent the variation of accretion and wind speeds in disk flows as functions of radial coordinate for j = 0.5. The solid line represents the accretion speed whereas the rest is for wind speed. It should be marked that the abosolute value of the velocities are been plotted here.

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Fig. 4(a) & 4(b) represent the variation of accretion and wind radial velocity (red) and sound speeds (green) in disk flows as functions of radial coordinate for j = 0.5, $\lambda_c = 2.4$, $\alpha_{SS} = 0.01$, $T = 7 \times 10^{11} \text{ K}(\text{ion temp})$ as we have assumed optically thin and geometrically thick flow.

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Solu	tions : $\frac{\eta}{s}$				



Fig. 4(c) & 4(d) represent the variation of accretion and wind $\frac{\lambda}{\lambda_k}$ and density (10⁻¹⁸gm/cc unit) in disk flows as functions of radial coordinate. All other parameters are been given the previous case.

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Solut	tions : $\frac{\eta}{s}$				



Fig. 4(e) represents the variation of accretion and wind $\frac{\eta}{s}$ for disk flows as functions of radial coordinate. $\frac{\eta}{s} \sim 4 \times 10^{-12}$ at the wind branch very near to the BH. The lower limit required 6×10^{-13} .

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Wind is dominating :

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Wind is dominating : DE in the form of MCG is prone to be thrown outwards from the accretion flows.



• Wind is dominating : DE in the form of MCG is prone to be thrown outwards from the accretion flows. As the negative pressure increases with the increase of *n*, the wind velocity becomes equal to the speed of light at a near vicinity of the the BH, when $\rho < 1$ which is generically true for flows around BHs.

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Conclusions and Future Scopes

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Conclusions and Future Scopes

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- We have a chance to reach the string theory lower bound if we count DE accretion. Setting the temperature will require to know the mean molecular weight of the flowing substance. Knowing which will merely lighten this problem.

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THANK YOU!

Ritabrata Biswas , IIEST Shibpur, Howrah. TCGC-2014 MCG Accretion : Limiting Value of η/s

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