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BH Thermodynamics with Cosmic Acceleration

Abhijit Mandal , IIEST Shibpur, Howrah. TCGC–2014

August 9, 2014

Black Hole Thermodynamics

Thermodynamics of Quintessence BHs 000000000

Conclusions

Effects of Cosmic Acceleration on Black Hole Thermodynamics

by Abhijit Mandal, Department of Mathematics, Indian institute of Engineering Science and Technology, Shibpur, Howrah-711 103, India. TCGC-2014

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- Black Hole as a Thermodynamic System
- Quintessence Black Holes
- Motive of the Current Work

2 Thermodynamics of Quintessence BHs

- Mass of QdS BHs
- Hawking Temperature
- Specific Heat
- Free Energy

3 Conclusions

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Conclusions

What a Black Hole is

 Highly dense collapsed massive star from which gravity prevents anything including light, from escaping.

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- Example : Xray binaries, Centre of galaxies, especially, active galaxies may host BHs. Andromeda galaxy M31 at its centre is hosting a super massive BH, OJ287 may be containing two supermasive BHs rotating each other etc.

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BH : Thermodynamic System

The laws of black hole mechanics are expressed in geometrized units.

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$$dE = rac{k}{8\pi} dA + \Omega dJ + \phi dQ$$

Where *E* is the energy, *k* is the surface gravity, *A* is the horizon area, Ω is the angular velocity, *J* is the angular momentum, ϕ is the electrostatic potential and *Q* is the electric charge.

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Since 1998 the cosmic acceleration hypothesis has been accepted by a large part of the astrophysicists and cosmologists followed by a series of observations¹.

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$$g_{tt}^{QdS} = 1 - \frac{r_g}{r} + \frac{Q^2}{r^2} - \frac{r^2}{a^2} - \left(\frac{r_q}{r}\right)^{3\omega_q + 1}$$

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Motive			

What is the destiny of the highly compact objects like BHs at "Big Rip"?

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Conclusions

Mass of Quintessence BHs

■ We get the form of the BH mass as a function of radius of event horizon(r_h), charge(Q), quintessence EoS(ω_q) and quintessence potential(r_q) as

$$M_{QdS} = \frac{r_h}{2} \left[1 + \frac{Q^2}{r_h^2} - \frac{r_h^2}{a^2} - \left(\frac{r_q}{r_h}\right)^{3\omega_q + 1} \right]$$



Fig.1b



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Conclusions

Mass of Quintessence BHs : Physical Interpretation

• The general trend of M_{QdS} vs r_h curve is primarily strictly decreasing and slowly increasing in the latter.

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Thermodynamics of Quintessence BHs

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- The general trend of M_{QdS} vs r_h curve is primarily strictly decreasing and slowly increasing in the latter.
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- According to Fig : 1(b) and 1(c), for $\omega_q = -0.66 \& -1$, Whenever the charge Q is nearly zero, the BHs are showing very small r_h and small masses.

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Outline

Black Hole Thermodynamics

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Hawking Temperature of QdS BHs

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Thermodynamics of Quintessence BHs

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$$T_{H} = \left(\frac{\partial M}{\partial S_{BH}}\right)_{Q} = \frac{1}{4r_{h}} - \frac{Q^{2}}{4r_{h}^{3}} - \frac{3r_{h}}{4a^{2}} + \frac{3\omega_{q}}{4r_{h}}\left(\frac{r_{q}}{r_{h}}\right)^{3\omega_{q}+1}$$

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Thermodynamics of Quintessence BHs

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Hawking Temperature of QdS BHs : Physical Interpretation

• At $\omega_q = -\frac{1}{3}$, the HT become negative for -2 < Q < 2.

Thermodynamics of Quintessence BHs

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- At $\omega_q = -\frac{1}{3}$, the HT become negative for -2 < Q < 2.
- In Fig : 2(a), whenever, $\omega_q = -0.4$ HT is positive if the charge is nearly zero.

Thermodynamics of Quintessence BHs

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- Now in Fig : 2(*b*) and 2(*c*), for $\omega_q = -0.66$ and -1 HT is positive for low charge.

Thermodynamics of Quintessence BHs

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Black Hole Thermodynamics

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Specific Heat of QdS BHs

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$$C_{Q,r_q} = T_H \left(\frac{\partial S_{BH}}{\partial T_H}\right)_{Q,r_q} = \frac{2r_h^2 \left[3r_h^4 + a^2 \left\{Q^2 - r_h \left(r_h + 3\omega_q r_q \left(\frac{r_q}{r_h}\right)^{3\omega_q}\right)\right\}\right]}{3r_h^4 + a^2 \left[r_h \left\{r_h + 3\omega_q (2 + 3\omega_q) r_q \left(\frac{r_q}{r_h}\right)^{3\omega_q}\right\} - 3Q^2\right]}.$$

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Fig.3a Fig.3b Fig.3c

Fig.3d

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Specific Heat of QdS BHs : Physical Interpretation

At the very first, for $\omega_q = -0.33$, for smaller r_h we get only negative C_{Q,r_q}

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The scenario for $\omega_q = -0.5$ is basically same with previous except the fact that with Q the increment of $r_{h,crit1}(Q)$ and decrease of $r_{h,crit2}(Q)$ is not as fast as the case we have analyzed before.

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Thermodynamics of Quintessence BHs

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Black Hole Thermodynamics

Thermodynamics of Quintessence BHs

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Conclusions

Free energy of QdS BHs

The expression for free energy becomes

Conclusions

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Free energy of QdS BHs

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$$F_{e} = \frac{r_{h}^{4} + a^{2} \left[3Q^{2} + r_{h} \left\{ r_{h} - (3\omega_{q} + 2)r_{q} \left(\frac{r_{q}}{r_{h}} \right)^{3\omega_{q}} \right\} \right]}{4a^{2}r_{h}}$$

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Fig.4a







Abhijit Mandal , IIEST Shibpur, Howrah. TCGC-2014 BH Thermodynamics with Cosmic Acceleration

Thermodynamics of Quintessence BHs

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Thermodynamics of Quintessence BHs

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Thermodynamics of Quintessence BHs ○○○○○○○●

Conclusions

Free energy of QdS BHs : Physical interpretation



Figures 5a-d represent the variation of free energy with respect to T_H for the different parameters set.

Fig 5a-5d show the variation of free energy with Hawking temperature. In all the cases except $\omega_q = -0.33$ we can find at least one cuspidal type double point where $\frac{\partial F_e}{\partial T_H}$ has single value but that occurs twice. This thing signifies a Hawking Page phase transition.

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Conclusio	ons		

■ So the obligatory sketch out of charged BH thermodynamics with the central engine amplified into a quintessence field is stable small BH → Second order phase transition → Unstable small/intermediate mass BH → First order phase transition → Stable intermediate mass BH → Second order phase transition → unstable super massive BHs.

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- So in a nutshell, before quintessence starts to work (\u03c6q = -0.33 > -\frac{1}{3}) it was preferable to have a small unstable BH followed by a large stable one. But in quintessence (-\frac{1}{3} > \u03c6q q > -1) BHs are destined to be unstable large one pre-quelled by stable/ unstable small/ intermediate mass BHs.

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Conclusions

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